Vector Calculus Without Vectors

Max Orchard

August 29, 2025

Vector Calculus Without Vectors (what MATH2901 could be, maybe)

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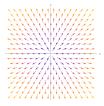
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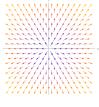


F(x,y) = (x,y)

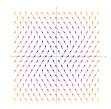
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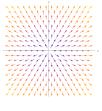


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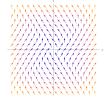
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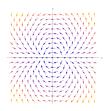
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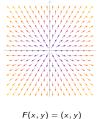
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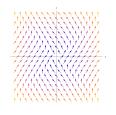
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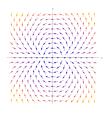
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We will denote the vector space of all vector fields on \mathbb{R}^n as $\mathfrak{X}(\mathbb{R}^n)$.

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Why are there multiple very different ways to differentiate a vector field?

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A 1-form is a covector field. That is, it is a (smooth) map $\omega: \mathbb{R}^n \to (\mathbb{R}^n \to \mathbb{R})$ that sends $p \in \mathbb{R}^n$ to a covector $\omega_p: \mathbb{R}^n \to \mathbb{R}$.



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We will denote the set of all 1-forms on \mathbb{R}^n as $\Omega^1(\mathbb{R}^n)$.



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Similarly, we denote by dy the 1-form that takes a vector to its y component:

$$dy_{(x,v)}(u,v)=v.$$



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 $\omega_{(x,y)}(u,v) = y^2u + 2x^2v$ is a 1-form (that is linear in u and v, but not in x and y). Here, $\omega = y^2 dx + 2x^2 dy$.



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Examples of 1-Forms

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In general, we cannot multiply 1-forms together and get another 1-form.

Structure of $\mathfrak{X}(\mathbb{R}^n)$ and $\Omega^1(\mathbb{R}^n)$

The spaces $\mathfrak{X}(\mathbb{R}^n)$ and $\Omega^1(\mathbb{R}^n)$ are both vector spaces (with pointwise addition and scalar multiplication):

$$(F+G)(\mathbf{x})=F(\mathbf{x})+G(\mathbf{x}), \quad (c\cdot F)(\mathbf{x})=c\cdot F(\mathbf{x}),$$

$$(\omega + \rho)_p(\mathbf{x}) = \omega_p(\mathbf{x}) + \rho_p(\mathbf{x}), \quad (c \cdot \omega)_p(\mathbf{x}) = c \cdot \omega_p(\mathbf{x}).$$

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Both vector spaces are *n*-dimensional, with bases

$$\{\partial x^1, \dots, \partial x^n\}$$
 for $\mathfrak{X}(\mathbb{R}^n)$, $\{\mathrm{d} x_1, \dots, \mathrm{d} x_n\}$ for $\Omega^1(\mathbb{R}^n)$,

where $\partial x^i(\mathbf{x}) = e^i = (0, \dots, \underbrace{1}_{i^{th} \text{ spot}}, \dots, 0).$



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where $\partial x^i(\mathbf{x}) = e^i = (0, \dots, \underbrace{1}_{i^{\mathsf{th}} \mathsf{spot}}, \dots, 0).$

Careful: the coefficients are functions $F_i(x_1, ..., x_n)$, not just scalars!

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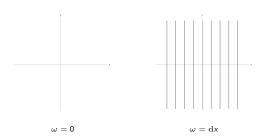




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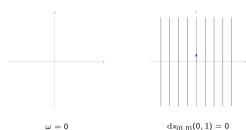
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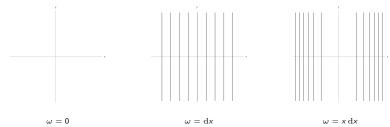


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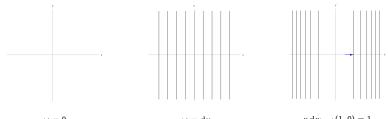
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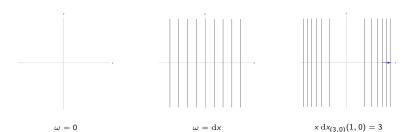
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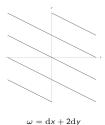
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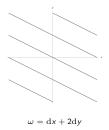
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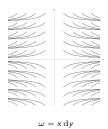
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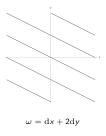


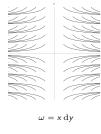


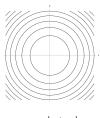






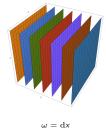






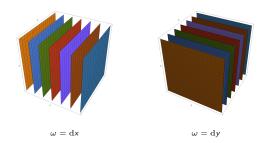
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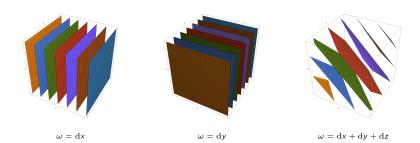


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The *flat operator* $\flat:\mathfrak{X}(\mathbb{R}^n)\to\Omega^1(\mathbb{R}^n)$ is defined by

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The musical isomorphisms give a one-to-one correspondence

$$\mathfrak{X}(\mathbb{R}^n) \longleftrightarrow \Omega^1(\mathbb{R}^n).$$



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In general, the following are identified under the musical isomorphisms:

$$\sum_{i=1}^n F_i(x_1,\ldots,x_n) \, \partial x^i \longleftrightarrow \sum_{i=1}^n F_i(x_1,\ldots,x_n) \, \mathrm{d} x_i.$$



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Example

Recall that $\partial x^i(\mathbf{x}) = e^i$. We have

$$(\partial x^i)_p^{\flat}(\mathbf{x}) = \langle e^i, \mathbf{x} \rangle = (\mathrm{d} x_i)_p(\mathbf{x}).$$

We have "lowered" the i, hence the name \flat .

In general, the following are identified under the musical isomorphisms:

$$\sum_{i=1}^n F_i(x_1,\ldots,x_n) \, \partial x^i \longleftrightarrow \sum_{i=1}^n F_i(x_1,\ldots,x_n) \, \mathrm{d} x_i.$$

For example, F(x, y) = (y, x) would be identified with y dx + x dy.



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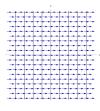


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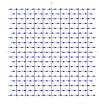
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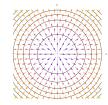
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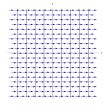


$$F = (x, y)$$

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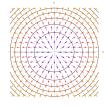
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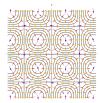
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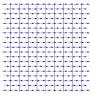
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This means we want to draw our lines *perpendicular* to the arrows in the vector field!

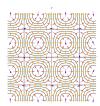






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The inverse operator \sharp can be visualised in a similar way.

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Definition

The *gradient* of a function $f: \mathbb{R}^n \to \mathbb{R}$ is the vector field

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This is a vector field, so can be turned into a 1-form through \flat . What does this look like?



Definition

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Remark: MATH1052/1072 says we have to instead dot with the unit vector $\hat{\mathbf{x}} = \mathbf{x}/\|\mathbf{x}\|$. We don't do this, otherwise $\mathrm{d}f$ wouldn't be linear.

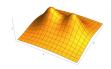


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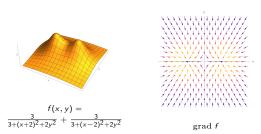
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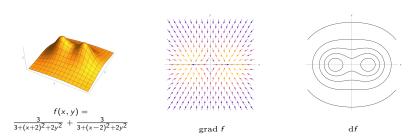
$$f(x,y) = \frac{3}{3+(y+2)^2+2y^2} + \frac{3}{3+(y-2)^2+2y^2}$$



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Enter: the 2-form.



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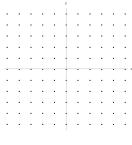
2-forms have two vector inputs, which draw a parallelogram.



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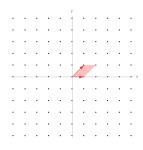


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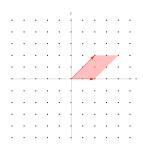
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$$dx \wedge dy_{(0,0)}((1,0),(1,1)) = 1$$



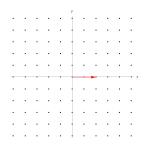
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$$dx \wedge dy_{(0,0)}((2,0),(2,2)) = 4$$



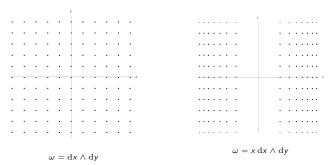
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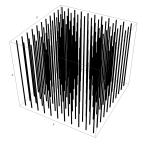
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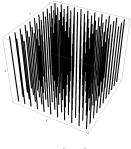


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Definition

The wedge product is the operator $\wedge: \Omega^1(\mathbb{R}^n) \times \Omega^1(\mathbb{R}^n) \to \Omega^2(\mathbb{R}^n)$ defined by

$$(\alpha \wedge \beta)(\mathbf{x}_1, \mathbf{x}_2) = \alpha(\mathbf{x}_1) \cdot \beta(\mathbf{x}_2) - \alpha(\mathbf{x}_2) \cdot \beta(\mathbf{x}_1).$$

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Note that the definition immediately implies $\alpha \wedge \alpha = 0$.



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It turns out that $\{dx_i \wedge dx_j\}_{i < j}$ is a basis for $\Omega^2(\mathbb{R}^n)$. That is, all 2-forms can be generated with just the wedge product alone.



Since the wedge product acts as multiplication, along each line in one direction, we need to count the number of times we hit lines in the other directions.



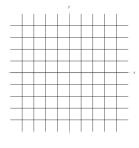
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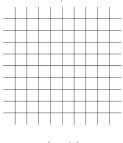


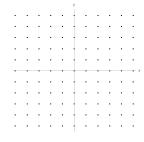
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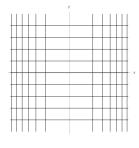




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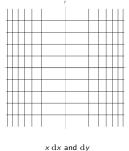


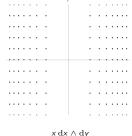
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The Exterior Derivative

We finally have the tools we need to define the derivative of a 1-form.



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Definition

The *exterior derivative* is the map $d: \Omega^1(\mathbb{R}^n) \to \Omega^2(\mathbb{R}^n)$ given on multiples of basis forms by

$$d(f dx_i) = \sum_{i=1}^n \frac{\partial f}{\partial x_i} dx_j \wedge dx_i$$

and extended additively.



Intuitively, if $\omega \in \Omega^1(\mathbb{R}^n)$, then $d\omega(\mathbf{x}, \mathbf{y})$ measures the *difference* in the change in $\omega(\mathbf{y})$ as you move along \mathbf{x} and the change in $\omega(\mathbf{x})$ as you move along \mathbf{y} :

$$d\omega(\mathbf{x}, \mathbf{y}) = \mathbf{x}(\omega(\mathbf{y})) - \mathbf{y}(\omega(\mathbf{x}))$$
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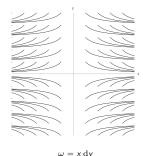


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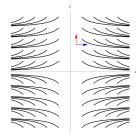
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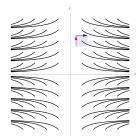
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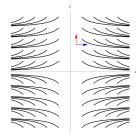
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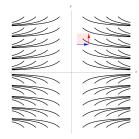
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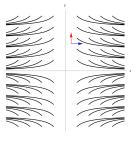
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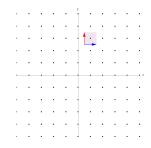




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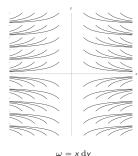


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The exterior derivative can therefore be seen as the *boundary* of the lines drawn by a form.

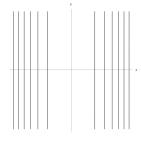
Vector Calculus Without Vectors





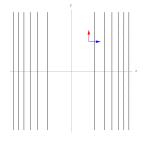
Intuitively, if $\omega \in \Omega^1(\mathbb{R}^n)$, then $d\omega(\mathbf{x}, \mathbf{y})$ measures the *difference* in the change in $\omega(\mathbf{y})$ as you move along \mathbf{x} and the change in $\omega(\mathbf{x})$ as you move along \mathbf{y} :

$$d\omega(\mathbf{x}, \mathbf{y}) = \mathbf{x}(\omega(\mathbf{y})) - \mathbf{y}(\omega(\mathbf{x}))$$
 (for constant \mathbf{x}, \mathbf{y}).



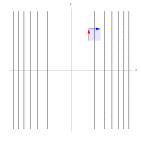
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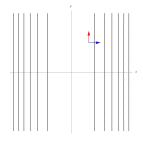
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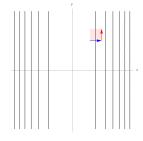






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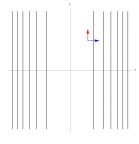


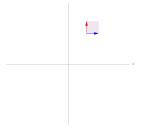




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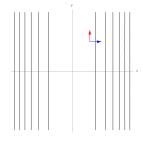




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The exterior derivative can therefore be seen as the *boundary* of the lines drawn by a form.



 $d\omega = 0$

 $\omega = x dx$

The above constructions can be generalised to k-forms (k-linear alternating maps $\underbrace{\mathbb{R}^n \times \cdots \times \mathbb{R}^n}_{k \text{ times}} \to \mathbb{R}$ assigned to each point).



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The wedge product and exterior derivative extend naturally, though ${\rm d}$ satisfies a "graded" product rule

$$d(\omega \wedge \alpha) = (d\omega) \wedge \alpha + (-1)^{k\ell} \omega \wedge (d\alpha),$$
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Important fact: by the symmetry of mixed partial derivatives, $d(d\omega) = 0$ for all ω . This corresponds to the fact that $\partial(\partial X) = \emptyset$.



As we've come to expect, a basis for $\Omega^k(\mathbb{R}^n)$ is given by

$$\left\{\mathrm{d} x_{i_1} \wedge \cdots \wedge \mathrm{d} x_{i_k}\right\}_{i_1 < \cdots < i_k}.$$



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If k > n, then $\Omega^k(\mathbb{R}^n) = \{0\}$. That is, $\Omega^n(\mathbb{R}^n)$ is the highest order possible. Moreover, $\dim(\Omega^n(\mathbb{R}^n)) = 1$.



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Remark: A scalar field $f: \mathbb{R}^n \to \mathbb{R}$ can be thought of as taking in zero vectors. Because of this, we say scalar fields are *0-forms*, and denote $C^{\infty}(\mathbb{R}^n,\mathbb{R})=\Omega^0(\mathbb{R}^n)$.



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In fact,
$$\star(\star\omega) = (-1)^{k(n-k)}\omega$$
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The Hodge star "completes" a form to the entire space.



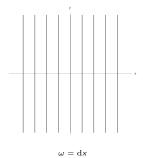
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The Hodge star can therefore be seen as the *orthogonal complement* of the lines drawn by a form.



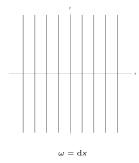
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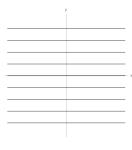
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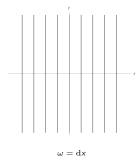


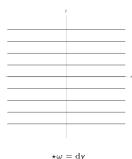


 $\star \omega = dv$

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The fact that $\star(\star\omega)=(-1)^{k(n-k)}\omega$ corresponds to the fact that $(U^{\perp})^{\perp} = U$ (and U = -U).



$$\mathfrak{X}(\mathbb{R}^3)$$



$$\mathfrak{X}(\mathbb{R}^3)$$

$$\Omega^1(\mathbb{R}^3)$$



$$\mathfrak{X}(\mathbb{R}^3)$$

$$\downarrow^{\flat}$$
 $\Omega^1(\mathbb{R}^3)$



$$\mathfrak{X}(\mathbb{R}^3)$$
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$$\sharp \dot{\mathbb{T}} \downarrow_{\flat}$$
 $\Omega^0(\mathbb{R}^3)$ $\Omega^1(\mathbb{R}^3)$



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$$\sharp igcap \downarrow_{\flat}$$
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$$\Omega^0(\mathbb{R}^3) \stackrel{\mathrm{d}}{\longrightarrow} \Omega^1(\mathbb{R}^3) \qquad \Omega^2(\mathbb{R}^3)$$

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This is the picture we've built up so far for \mathbb{R}^3 :

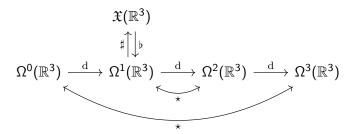
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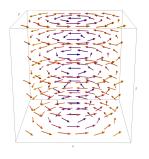
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Let's do some calculus! (Finally!)

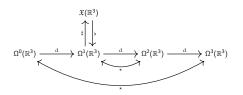


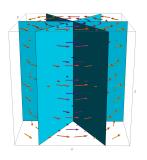
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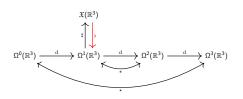


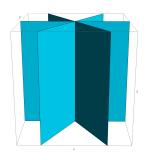
$$F = (-y, x, 0)$$



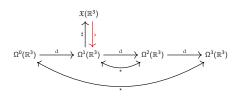


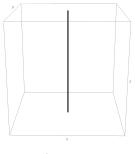
$$F^{\flat} = -v \, \mathrm{d}x + x \, \mathrm{d}v$$



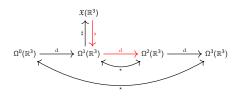


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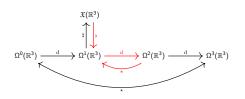


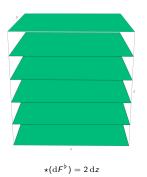


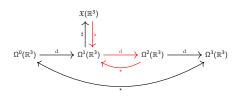


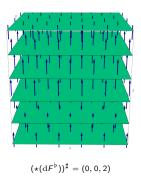


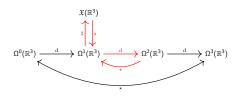


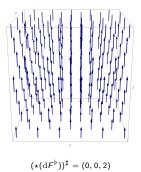


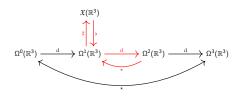










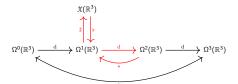


Curl as an Exterior Derivative

Definition

The *curl* is the operator $\operatorname{curl}:\mathfrak{X}(\mathbb{R}^3)\to\mathfrak{X}(\mathbb{R}^3)$ defined by

$$\operatorname{curl}(F) = (\star(\mathrm{d}F^{\flat}))^{\sharp}.$$

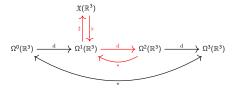


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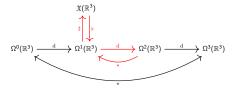


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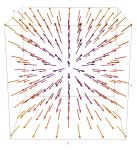


The \star , \flat , \sharp kind of obscure what's going on: they are just isomorphisms allowing us to identify one space with another. What we're really doing is differentiating a 1-form and interpreting it as a vector field.

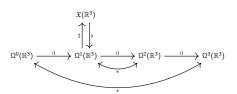
In a similar way, we can see the divergence as the *derivative of a 2-form*.

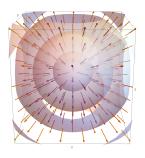




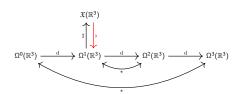






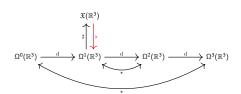


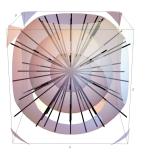
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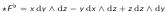


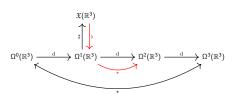


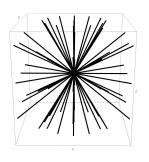
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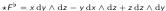


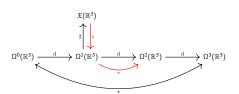


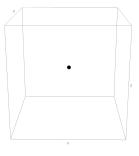




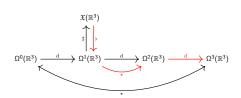




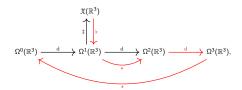




$$d(\star F^{\flat}) = 3 dx \wedge dy \wedge dz$$



In a similar way, we can see the divergence as the *derivative of a 2-form*. Let's consider F(x, y, z) = (x, y, z).

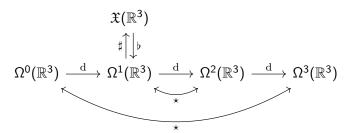


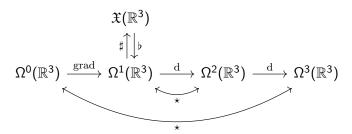
Definition

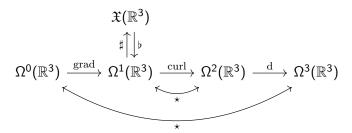
The *divergence* is the operator $\mathrm{div}:\mathfrak{X}(\mathbb{R}^3)\to C^\infty(\mathbb{R}^3,\mathbb{R})$ defined by

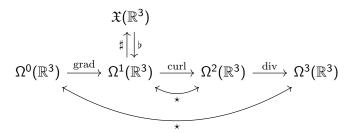
$$\operatorname{div}(F) = \star \operatorname{d}(\star F^{\flat}).$$

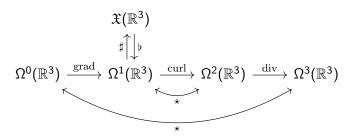










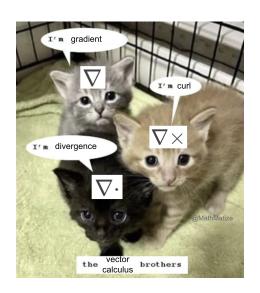


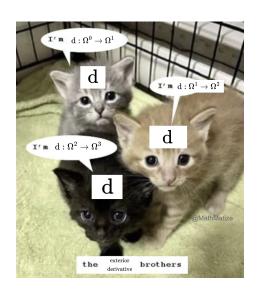
Exercise

We know (MATH2001) that conservative vector fields have zero curl:

$$\nabla \times (\nabla f) = 0.$$

Prove this fact in two lines using the framework of forms. See if you can come up with another similar fact.





Integrating Forms

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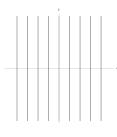
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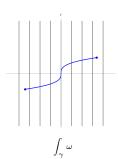
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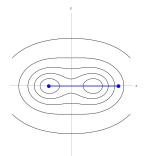
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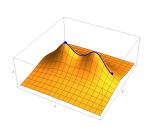


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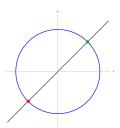
This is actually saying something very trivial: the number of lines that enter (and don't exit) a region is the number of lines that enter (and don't exit) a region.

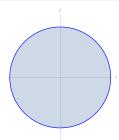
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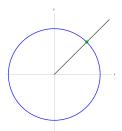


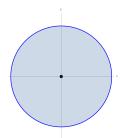
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By varying U and ω , we recover many classical vector calculus theorems.

| U | ∂U | ω | $\mathrm{d}\omega$ | Theorem |
|-----------------|---------------------------|---------------------|----------------------------------|--------------------|
| [a, b] | { a, b} | f | f'(x) dx | FTC |
| $\gamma([a,b])$ | $\{\gamma(a),\gamma(b)\}$ | f | $(\nabla f)^{\flat}$ | FTLI |
| U | ∂U | $\omega\in\Omega^1$ | $d\omega$ | Green's theorem |
| S | ∂S | F♭ | $\operatorname{curl}(F)^{\flat}$ | Stokes' theorem |
| V | ∂V | ∗F♭ | $\star \operatorname{div}(F)$ | Divergence theorem |

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- connections and the Yang-Mills equations.



References

For some further reading:

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